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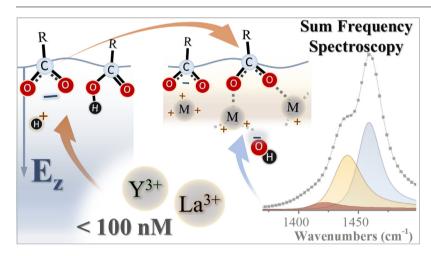
# La<sup>3+</sup> and Y<sup>3+</sup> interactions with the carboxylic acid moiety at the liquid/vapor interface: Identification of binding complexes, charge reversal, and detection limits



Adrien Sthoer <sup>a</sup>, Ellen M. Adams <sup>b,d</sup>, Sanghamitra Sengupta <sup>a,e</sup>, Robert W. Corkery <sup>a,c</sup>, Heather C. Allen <sup>b</sup>, Eric C. Tyrode <sup>a,\*</sup>

- <sup>a</sup> Department of Chemistry, KTH Royal Institute of Technology, SE-10044 Stockholm, Sweden
- <sup>b</sup> Department of Chemistry & Biochemistry, The Ohio State University, Columbus, OH 43210, USA
- Department of Applied Mathematics, Research School of Physics and Engineering, Australian National University, Canberra, ACT0200, Australia
- <sup>d</sup> Department of Physical Chemistry II, Ruhr-Universität Bochum, 44801 Bochum, Germany<sup>1</sup>
- e Ultrafast Spectroscopy, AMOLF, 1098 XG Science Park, Amsterdam, The Netherlands

#### G R A P H I C A L A B S T R A C T



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### ABSTRACT

Specific interactions of yttrium and lanthanum ions with a fatty acid Langmuir monolayer were investigated using vibrational sum frequency spectroscopy. The trivalent ions were shown to interact with the charged form of the carboxylic acid group from nanomolar concentrations (<300 nM). Analysis of the spectral features from both the symmetric and the asymmetric carboxylate modes reveals the presence of at least three distinct coordination structures linked to specific binding configurations. Although the same species were identified for both La³+ and Y³+, they display a different concentration dependence, highlighting the ion-specificity of the interaction. From the analysis of the response of interfacial water molecules, the reversal of the surface charge, as well as the formation of yttrium hydroxide complexes, were detected upon increasing the amount of salt in solution. The binding interaction and kinetics of absorption are sensitive to the solution pH, showing a distinct ion speciation in the interfacial region

<sup>\*</sup> Corresponding author.

E-mail address: tyrode@kth.se (E.C. Tyrode).

<sup>&</sup>lt;sup>1</sup> Authors present addresses.

Vibrational sum frequency spectroscopy Molecular characterization Carboxylate binding complexes Charge reversal Poisson Boltzmann theory Vibrational spectroscopy Arachidic acid Eicosanoic acid monolayer when compared to the bulk. Changing the subphase pH or adding a monovalent background electrolyte that promotes deprotonation of the carboxylic acid headgroup could further improve the detection limit of  $La^{3+}$  and  $Y^{3+}$  to concentrations < 100 nM. These findings demonstrate that nM concentrations of trace metals contaminants, typically found on monovalent salts, can significantly influence the binding structure and kinetics in Langmuir monolayers.

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#### 1. Introduction

The binding of transition metal ions to charged surfaces is vital in many processes, such as protein folding,[1] atmospheric chemical reactions,[2] and geochemical sequestration and transport of contaminants already from sub-micromolar concentrations,[3] A multitude of studies have demonstrated the ion specific nature of surface binding phenomena, in which ions with the same charge can interact and subsequently alter monolayer properties in different manners,[4–11] Long chain fatty acids have been used as a model for negatively charged surfaces to investigate the cation binding due to the prevalence of carboxylic acid groups in biological and atmospheric chemistry systems,[12–14]

Grazing incidence X-ray diffraction measurements and vibrational spectroscopic studies of metallic soaps have been combined in an attempt to correlate the coordination structures of the carboxylate moiety with its vibrational signatures and the ion-specific nature of the metallic cations.[15–17] Similar investigations have been carried out for hydrated carboxylate anions in aqueous solution.[18] However, less is known about coordination structures of metallic cations to carboxylates constrained at interfaces, especially for polyvalent ions. Existing studies of carboxylate interfaces report contrasting results and interpretations for the type of binding interactions.[19–23] This is especially true in the case of monovalent and divalent salts, where strong discrepancies were observed in the vibration frequency of the metal-carboxylate complex, which could vary with time, even for highly concentrated solutions.[22,24,25]

While differences in experimental parameters, such as pH, could account for some of the observed discrepancies, a recent study by Sthoer et al. has revealed that trace amounts of polyvalent ions in solution can compete with the ion of interest for adsorption sites and cause variations in the vibrational pattern of the fatty acid monolayer.[26] While the impact of organic impurities included in the salts used had been previously demonstrated, [24,27] the latter study showed that trace amounts of polyvalent metal impurities, such as trivalent ions, must also be considered when investigating binding to a charged interface.[26]

Studies of trivalent cation binding are challenging due to the limitations in predicting their speciation in aqueous solutions and their complex electronic structures. Nevertheless, trivalent ions are known to play a central role in protein and colloidal reactivity, triggering crystallization , phase separation, and surface charge reversal.[28–30] For example, the adsorption of Fe<sup>3+</sup> and La<sup>3+</sup> on bacterial and bovine serum albumin proteins was revealed to be strongly pH dependent at micromolar concentrations,[31,32] and caused a charge reversal of lipid surfaces starting from submicromolar LaCl<sub>3</sub> concentrations,[33,34]

The complexation of the trivalent ions La<sup>3+</sup> and Fe<sup>3+</sup> to fatty acid Langmuir monolayers has been probed by X-ray reflectivity, X-ray fluorescence, and more recently vibrational sum frequency spectroscopy.[35–38] Specific ion effects were observed, in which it was concluded that La<sup>3+</sup> adsorption to the carboxylic acid was led only by Coulombic interactions. On the other hand, the interaction of Fe<sup>3+</sup> was determined to be of covalent nature in a pH range where iron hydroxide complexes are absent from solution.[35,36]

Vibrational sum frequency investigations using millimolar salt concentrations confirmed the ion specificity and pH dependence of the ionic interactions by probing the response of water molecules in the interfacial region; however, little attention was brought to the direct molecular interactions with the carboxylic acid headgroup, the charge reversal and co-ion binding.[38] To the authors' knowledge, no molecular level observations have been reported for Y³+ ion interactions with fatty acid Langmuir monolayers.

In this study vibrational sum frequency spectroscopy is used to probe the complexation of trivalent ions to an arachidic acid Langmuir monolayer. La<sup>3+</sup> and Y<sup>3+</sup> were chosen as they both usually have the same oxidation state in solution (i.e. + 3) and possess an accommodating speciation diagram in the pH range where the carboxylic acid headgroup deprotonates. Moreover, they both interact through their s electrons' shell, which gives a spherical symmetry and considerable flexibility to form similar coordinated complexes. This provides an opportunity to independently investigate the effect of ion size on the type of coordination formed with the carboxylate. Yttrium shares many similarities with the lanthanides, particularly gadolinium, terbium, dysprosium, holmium, and erbium that, due to the lanthanide contraction, have comparable ionic radii. Additionally, La<sup>3+</sup> and Y<sup>3+</sup> are of interest as they are frequently used in microelectronics and the newest hybrid materials[39] and are typically present as impurities in monovalent and divalent salts. We find that the binding interaction is concentration dependent and specific to each cation. The deprotonation of the carboxylic acid at natural pH begins at concentrations as low as 300 nM. This detection limit is further decreased to 100 nM by adding a NaCl background electrolyte, which promotes the binding of Y<sup>3+</sup>. The interfacial water response is also systematically compared with the carboxylate binding pattern. Water bands prove to be more sensitive to the first stages of deprotonation, as well as the charge reversal of the monolayer at higher ionic strengths. Furthermore, ion adsorption kinetics are found to be impacted by changes in the subphase pH. Taken together, these results demonstrate that trivalent ions have specific spectroscopic signatures that could be used to identify trace metal contaminants in solution and that ion speciation as a function of pH and adsorption kinetics must be considered in diffusion-limited cation binding investigations.

### 2. Materials and methods

### 2.1. Materials

YCl<sub>3</sub> (99.99%, trace metals basis, anhydrous), LaCl<sub>3</sub>.7H<sub>2</sub>O (99.999%, trace metals basis), NaCl (99.999%, trace metals basis), NaOH (99.99%, trace metals basis), deuterated arachidic acid (eicosanoic-d39 acid, 97%), arachidic acid (eicosanoic acid, 99%), and chloroform (anhydrous grade, stabilized with ethanol) were obtained from Merck. HCl 36.5% (99.999%, trace metals basis) was purchased from Alfa-Aesar. Ultrapure water was obtained from an Integral 15 Millipore system with a resistivity of 18.2 MΩ.cm and a total organic content < 3 ppb. NaCl was baked at 500 °C for 1 h to remove potential traces of organic contami-

nants. All other compounds were used as received. The glassware was cleaned by a three-step sonication procedure, first in ethanol, followed by an alkaline cleaning agent solution (Deconex from Borer Chemie), and finally, ultrapure water. The glassware was thoroughly rinsed in ultrapure water between each step. In order to avoid ionic specific interactions at the surface of the pH sensor glass electrode, the pH was adjusted prior to the addition of salt and later cross-checked with multiple pH paper indicators (MColorpHast<sup>TM</sup> pH 0 – 14, 4.0 – 7.0, 6.5 – 10, and 11.0 – 13.0 and Whatman® Panpeha<sup>TM</sup> pH 0 – 14 from Merck). To test the potential effect of bulk depletion due to ion adsorption to the glass surfaces at the lowest salt concentration measured, solutions were also prepared using polypropylene volumetric flasks. However, no measurable differences could be detected, suggesting that the effect is negligible.

### 2.2. Fatty acid compression isotherms

The compression isotherms and the sum frequency measurements on the fatty acid monolayers were carried out in a temperature-controlled trough (KSV NIMA, 50 mm width, 9,750 mm<sup>2</sup>) made out of Teflon<sup>TM</sup> with Delrin<sup>TM</sup> compression barriers. The surface pressure was measured with a 10 mm wide Wilhelmy paper plate. Fatty acid depositing solutions were prepared in chloroform. The spectroscopy measurements were performed no less than 10 min after the deposition of the monolayer to ensure the evaporation of the solvent. The temperature of the subphase was kept constant at 22 °C ± 0.1 °C. The trough was carefully cleaned before every experiment and soaked overnight with a saturated solution of EDTA to remove any potential traces of trivalent ions from previous experiments ( $\sim$ 3 mM at 25 °C). The EDTA procedure was found to be an important step to ensure reproducible results, particularly at sub µM concentrations, as evidence was found that yttrium and lanthanum cations could adsorb (or maybe even absorb) to the Teflon trough or the glass window located in the middle of the trough. All the sum frequency spectroscopy measurements presented in this article were performed at a constant surface pressure of 20 mN/m. At this pressure, the average area per molecule in the monolayer is approximately equal for all sub-phase salt concentrations ( $\sim$ 19 Å<sup>2</sup>/molecule) measured, which simplifies the spectral analysis.

# 2.3. Vibrational sum frequency spectroscopy

The vibrational sum frequency spectrometer has been described in detail elsewhere. [40] Briefly, a fraction of the output from a Ti: Sapphire amplifier ( $\sim$ 90 fs, 1 kHz,  $\sim$ 6 W, Integra-C, Amplitude Technologies, France) is used to pump a HE-Topas (Light Conversion, Lithuania), to obtain a broadband tunable IR pulse. The remaining fraction of the amplifier output is sent to a home-built pulse shaper, to generate a bandwidth tuneable 805 nm picosecond pulse. The visible and IR beams are focused on the sample position in a co-propagating geometry at angles of incidence of 70° and 55°, and average powers of  $\sim$  30 mW and  $\sim$  4 mW, respectively. The setup displays a high degree of automation, which allows collecting spectra in wide spectral regions. The spectral resolution in these experiments was set to < 3 cm<sup>-1</sup> by closing the slit in the beam shaper. Spectra were recorded in three polarization combinations, mainly SSP, SPS, and PPP, and normalized by the non-resonant sum frequency response from a gold surface.[40] To quantify the deprotonation of the monolayer, as well as the respective contributions of the metal-carboxylate complexes, the spectra were fitted with a convolution of Lorentzian and Gaussian line shapes, which account respectively for homogeneous and inhomogeneous broadening (see SI for additional details).[41,42] All the sum frequency spectra are presented in normalized units

(n.u.) using as reference the  $r^+$  mode in the SSP spectrum of arachidic acid on a pure water subphase. The IR laser path up to the sample position is purged with an overpressure of filtered,  $CO_2$ -free dry air, which helps keep the subphase pH constant within 0.2 pH units during the time of measurements (up to 1 h). This is particularly relevant for experiments performed at a pH > 7.

#### 3. Results and discussion

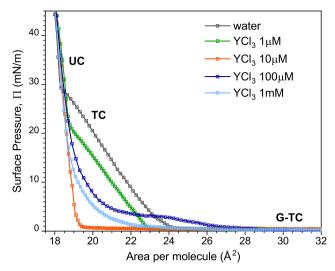
3.1. Surface pressure ( $\Pi$ ) vs. Molecular area (A) isotherms: Effect of  $Y^{3+}$  and  $La^{3+}$ 

Langmuir isotherms of arachidic acid on YCl3 and LaCl3 solutions were recorded for concentrations varying from 100 nM up to 1 mM. Fig. 1 shows the Π-A isotherms for four selected concentrations of YCl<sub>3</sub>, as well as that for a pure water subphase. For concentrations below 1 µM YCl<sub>3</sub> the isotherms are largely similar to that on pure water, displaying three distinct regions: a coexistence region between a 2D gas (G) and a tilted condensed (TC) phase at molecular areas >  $\sim$ 24 Å<sup>2</sup>, a purely TC phase, and an untilted condensed (UC) phase at surface pressures above  $\sim 27$  mN/m.[43–46] At 1 μM YCl<sub>3</sub> changes become apparent in the isotherm, with the TC-UC phase transition pressure decreasing to ~ 20 mN/m. At 10 μM the TC phase is no longer present, leaving the 2D gas phase in direct equilibrium with a UC phase, in agreement with a previously reported isotherm for this specific concentration.[47] Further increasing the salt concentration leads to a small expansion of the monolayer with the appearance of a non-horizontal plateau starting from a surface pressure of  $\pi_t \approx 3$  mN/m at 100  $\mu$ M, and the replacement of the 2D gas coexistence region by a liquid expanded (LE) phase. Finally, at 1 mM YCl<sub>3</sub> no discontinuities are apparent in the isotherm, showing a continuous transition between the LE and UC phase at the highest surface pressures. The  $\pi$ -A isotherms of arachidic acid on LaCl<sub>3</sub> solutions show a very similar concentration dependent behaviour (see SI).

The surface equilibria described by the Langmuir isotherms reflect a subtle balance between cohesive and repulsive forces. The former include short range Van der Waals interactions between the alkyl chains, headgroup bridging through hydrogenbonding, and complex formation, while the latter refers to short and long range forces opposing the packing, such as steric overlap, electrostatic repulsion between charged headgroups, as well as diffusion and thermal motion.[43,48-51] The contraction of the monolayer to the densest UC phase reflect an increase of the cohesive forces, as observed, for example, in stearic acid monolayers interacting with divalent ions at high concentrations.[52] On the contrary, monolayer expansion and appearance of the plateau of phase transition with a LE phase, result from the introduction of new repulsive forces, which in the case of fatty acid monolayers on monovalent chloride salt solutions originate from the charging of the monolayer. [53,54] The Langmuir isotherms recorded for YCl<sub>3</sub> and LaCl<sub>3</sub> indicate that the interfacial equilibrium is affected in both ways, revealing a complex interaction between the trivalent ion salts and the carboxylic acid headgroup. A molecular level insight into these interactions can be obtained using vibrational sum frequency spectroscopy.

# 3.2. $Y^{3+}$ and $La^{3+}$ interactions with the carboxylate headgroup: Concentration dependence

The sum frequency spectra of deuterated arachidic acid Langmuir monolayers on solutions of varying concentrations of YCl<sub>3</sub> and LaCl<sub>3</sub> under the SSP polarization combination are shown in Fig. 2a and 2b. The spectral region shown targets vibrational modes from the deuterated alkyl chain and the carboxylic acid headgroup.



**Fig. 1.** Surface pressure vs. molecular area isotherms of an arachidic acid monolayer on a pure water subphase and different YCl $_3$  concentrations. pH  $\sim$  5.8 and temperature 22 °C. The phases in equilibrium depicted (UC, TC, and G) correspond to those of the pure water subphase. The isotherms on LaCl $_3$  solutions show a very similar concentration dependent behaviour and can be found in the SI.

The deuteration of the chain facilitates the interpretation of the spectra as the headgroup vibrations no longer overlap with the methyl and methylene deformation bands. [46] In the CD stretching region, the three sharp peaks centered at  $\sim 2075~{\rm cm}^{-1},$   $\sim\!2135~{\rm cm}^{-1},$  and  $\sim 2218~{\rm cm}^{-1},$  are all linked to the terminal methyl group and assigned to the symmetric (r+), Fermi resonance (r+R), and asymmetric stretches (r-), respectively. [46,55] The lack of features associated with the CD2 groups indicate that the alkyl chain has a high degree of conformational order, consistent with the high packing density in the monolayer in an all-trans configuration. The vibrational modes from the chain appear insensitive to increasing amounts of trivalent ions in solution, and the phase transition between the tilted to untilted condensed phases (TC  $\rightarrow$  UC). This is rationalized by the similar area per molecule expected at 20 mN/m for all salt concentrations (see Fig. 1).

In contrast, significant changes are observed in the headgroup vibrations. Up to  $\sim 100$  nM, the spectra look similar to that of pure water, where only features from the uncharged carboxylic acid headgroup are detected, mainly the broad C-OH stretch at  $\sim 1300~\text{cm}^{-1}$ , and the carbonyl stretch C=O at  $\sim 1720~\text{cm}^{-1}$  (see inset in Fig. 2a).[46] However, starting from 300 nM, the deprotonation of the monolayer becomes apparent for both YCl<sub>3</sub> and LaCl<sub>3</sub> with the appearance of carboxylate stretching modes between 1400 and 1500 cm $^{-1}$ . This sets the detection limits for Y $^{3+}$  and La $^{3+}$  in the range between 100 and 300 nM.

Further increasing the salt concentration in the subphase leads to multiple vibrational contributions in the carboxylate region, as well as the decrease of the C=O mode from the uncharged carboxylic acid. Three resolved features can be readily identified in the symmetric carboxylate stretching region. They are centered at approximately  $\sim 1416~{\rm cm^{-1}}$  (Peak 1, ionic),  $\sim \! \! 1438~{\rm cm^{-1}}$  (Peak 2, bridging bidentate), and  $\sim 1458~{\rm cm^{-1}}$  (Peak 3, chelating bidentate), as shown in the inset in Fig. 2b. The same modes are observed for lanthanum and yttrium, but their relative intensities and concentration dependence vary depending on the identity of the cation. Though the assignments will be discussed in more detail in the next section, it is assumed that each symmetric carboxylate stretching contribution corresponds to a defined ion-pair or binding configuration between the metallic cation and the carboxylate.

The relative changes are better visualized in Fig. 2c and 2d, where the fitted amplitudes for the three symmetric carboxylate

peaks, and the carbonyl stretch, are shown as a function of ion concentration. The figures summarize the results from multiple repeat experiments (see SI for details of the fitting procedure). The C=O signal for both salts starts decreasing at sub µM concentrations and is no longer present at 10 µM, indicating that the monolayers are fully deprotonated beyond this point (the apparent binding constants for pK<sub>Y</sub> and pK<sub>La</sub> are both 6.0  $\pm$  0.1). This is in stark contrast to the predictions of the Gouy Chapman model for 1:3 salts, where the monolayer is expected to be  $\sim 10\%$  deprotonated at 10  $\mu$ M and just  $\sim$  25% at 1 mM (see SI). This adds support for strong specific interactions between the trivalent cations and the carboxylic acid group, which are not accounted for in Poisson-Boltzmann theory. Moreover, the absence of the hydrated carboxylate symmetric stretch at  $\sim 1408 \ \text{cm}^{-1}\text{,}$  which has been observed in fatty acid monolayers on solutions containing monovalent ions that do not specifically interact with the carboxylate. [26,46] show that complex formation occurs from the lowest trivalent ion concentrations.

Focusing on the amplitudes of the symmetric carboxylate peaks, the trends for the yttrium cation show a typical Langmuir behaviour for all the three contributions, where saturation levels are reached at  $\sim 5~\mu M$  (Fig. 2c). On the contrary, for the bulkier lanthanum cation, a steady-state in the binding configuration is only reached at 10 mM, with two carboxylate contributions (*Peak 1* at  $\sim 1416~cm^{-1}$  and *Peak 3* at  $\sim 1458~cm^{-1}$ ) showing opposite trends once the monolayer is fully deprotonated (Fig. 2d). These contrasting behaviours demonstrate differences in the binding mechanisms as well as the binding-pattern proportions between the two cations. Given that the cross-sections for each species are not known with sufficient precision, the absolute proportions between the different complexes in the monolayer cannot be accurately determined. Nonetheless, an estimate is proposed in the SI.

# 3.3. Interfacial binding patterns: Assignment of the carboxylate stretching modes

The identification of the type of coordination from the vibrational frequencies of the carboxylate stretches relies on previous correlations established by comparison with crystallographic structures in solid-state acetate crystals,[15,56] and in aqueous acetate solutions.[18] The relevant parameters to consider are the centre positions of the symmetric and asymmetric stretching modes ( $v_{coo-,as}$ ,  $v_{coo-,as}$ ), and their relative difference in frequencies  $\Delta(v_{as} - v_s)$ . However, it has been recognized that additional factors could influence the vibrational pattern observed. They include the polarity of the metal-oxygen bonds, the identity of the cation, the geometry of the O—C—O angles, or even the number of water molecules interacting with the charged species.[15,18,57] Despite these limitations, a general classification, albeit with many exceptions, has been proposed for  $\Delta(v_{as} - v_s)$ , where  $\Delta_{Unidentate} > \Delta_{Ionic} > \Delta_{Bridg}$  $_{ing/Chelating}$  with  $\Delta_{Chelating}$  (<90 cm<sup>-1</sup>) often <  $\Delta_{Bridging}$ .[16,58] The typical coordination possibilities between metal cations and the carboxylate group are depicted in Fig. 3. Except for the ionic / ion-pair form that results from electrostatic interactions,[59] the metal-carboxylate coordinative binding complexes are commonly sorted into four distinct classes, namely unidentate (I), chelating bidentate (II), bridging bidentate (III), and polymeric/tridentate forms (IV).[15,17,56] The polymeric form is a mixture between the chelating and bridging bidentate, where at least one of the oxygen atoms is simultaneously coordinated to two metallic cations, which in turn can also be involved in a type (II) or type (III) binding with another carboxylate group.[15,17]

As shown in the previous section, in the sum frequency spectra, the symmetric stretches are predominantly observed under the SSP polarization combination. The asymmetric carboxylate stretches are, in turn, mainly detected in the SPS polarization, with

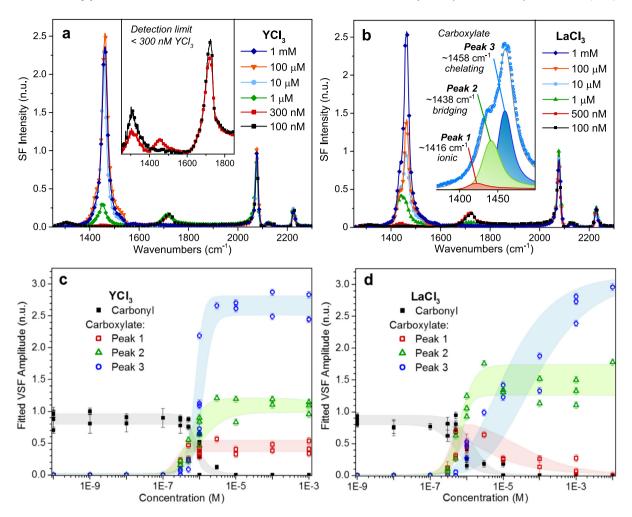


Fig. 2. Vibrational sum frequency spectra of a deuterated arachidic acid Langmuir monolayer on an aqueous subphase containing 100 nM up to 1 mM of a) YCl<sub>3</sub> and b) LaCl<sub>3</sub> at pH  $\sim$  5.8. Spectra were recorded in the SSP polarization combination. The surface pressure was set to 20 mN/m and the temperature to 22 °C. In the inset in b) three contributions can be distinguished for the carboxylate symmetric modes that are associated with specific metal-carboxylate coordination structures: ionic (Peak 1), bridging bidentate (Peak 2), and chelating bidentate (Peak 3)(assignment are discussed in a later section). c) Fitted amplitudes for the three carboxylate peaks and carbonyl stretch as a function of YCl<sub>3</sub> and d) LaCl<sub>3</sub> concentration for multiple repeat experiments (error bars reported are linked to the fitting procedure). Shaded regions are added as guides to the eyes to distinguish the trends for the different peaks. Note that in contrast to the data collected at higher pH (Fig. 8), the spectral features at pH  $\sim$  5.8 showed no time dependence in any of the concentrations measured.

PPP spectra showing contributions from both. Spectra recorded in the three polarization combinations for deuterated arachidic acid monolayers on 1 mM YCl<sub>3</sub> and LaCl<sub>3</sub> subphases are presented in Fig. 4. Besides the symmetric carboxylate stretches previously described for the SSP spectra, in the SPS spectra, several resolved features are observed at higher frequencies. Assigned to asymmetric stretches, they are centered at  $\sim$  1520 cm $^{-1}$  (Peak 4),  $\sim$ 1540 cm $^{-1}$  (Peak 5), and  $\sim$  1575 cm $^{-1}$  (Peak 6). In the inset of Fig. 4b, where an enlarged view of the SPS and PPP spectra is presented for LaCl<sub>3</sub>, the contributions from the three symmetric and the three asymmetric modes can be clearly resolved. The average peak positions for the six bands obtained from the fits of all spectra collected are summarized in Table 1.

Each symmetric carboxylate mode should have an associated asymmetric band; however, the pairing is not entirely obvious given the number of features observed. The assignment of *Peak 1* at  $\sim 1416~\rm cm^{-1}$  is the most straightforward, as it lies in the frequency range of ion pairing interactions (solvent-shared /contact ion pairs) with the carboxylate headgroup.[26,54] Thus, in agreement with Marcus et al.'s classification, it arises from pure electrostatic interactions and likely associated with a solvent separated ion pair with the hydrated trivalent ion.[59] Based on the same

studies,[26,54] the accompanying asymmetric stretch for *Peak 1* is expected between 1535 and 1540 cm<sup>-1</sup>, but with an order of magnitude lower intensity. Consequently, given the relatively weak amplitude of *Peak 1* (Fig. 2c and 2d), its asymmetric stretch should be barely detectable and, if anything, contributing to the intensity of *Peak 5*.

The asymmetric stretches linked to *Peak 2* at  $\sim 1438 \text{ cm}^{-1}$ , and Peak 3 at  $\sim 1458 \text{ cm}^{-1}$ , can be determined by comparing the evolution of their fitted amplitudes as a function of salt concentration, particularly in the case of LaCl<sub>3</sub> where the two bands are more readily resolved. It is concluded that Peak 2 is associated with Peak 5 at  $\sim 1545~\text{cm}^{-1}$  and Peak 3 with Peak 4 at  $\sim 1525~\text{cm}^{-1}$  (see SI for an extended description of this procedure). The Peak 2-Peak 5 symmetric/asymmetric correlation is also further supported by previous studies on Mn<sup>2+</sup> and Ni<sup>2+</sup> interactions with carboxylic acid monolayers, where single bands linked to the symmetric and asymmetric stretches were observed at equivalent peak positions. Nonetheless, the symmetric pair for the weak asymmetric contribution at  $\sim 1577 \text{ cm}^{-1}$  (*Peak 6*) remains uncertain (possibly linked to the polymeric complexes). The proposed assignments for both trivalent salts are summarized in Table 2. The chelating bidentate coordination (II) (Peak 3 - Peak 4) exhibiting a particularly small

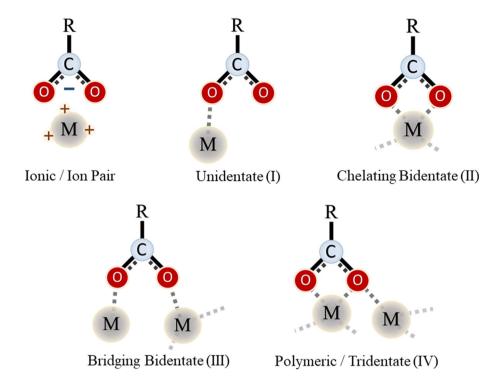


Fig. 3. Schematic representation of the typical coordination modes for the carboxylate group with metallic cations.

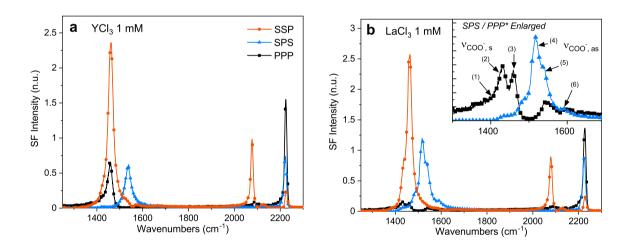


Fig. 4. Vibrational sum frequency spectra recorded under the SSP, SPS and PPP polarization combinations of deuterated arachidic acid monolayers on a) 1 mM YCl<sub>3</sub> and b) 1 mM LaCl<sub>3</sub> solution at pH  $\sim$  5.8 and a constant surface pressure of 20mN/m. The inset b) is an enlargement of the SPS and PPP spectra, where the PPP\* spectrum has been multiplied by 5 for ease of comparison. Spectra were recorded in triplicate, showing all the same trends.

**Table 1**Vibrational frequencies of the symmetric  $v_{coo-,s}$  and asymmetric  $v_{coo-,as}$  stretches of the carboxylate headgroup of the arachidic acid monolayers on YCl<sub>3</sub> and LaCl<sub>3</sub> subphases.

	Symmetric v <sub>coo-,s</sub>			Asymmetric $v_{coo-,as}$		
	Peak 1 (ionic)	Peak 2 (bridging)	Peak 3 (chelating)	Peak 4 (chelating)	Peak 5 (bridging)	Peak 6
YCl <sub>3</sub> LaCl <sub>3</sub>	1416 ± 1 cm <sup>-1</sup> 1416 ± 1 cm <sup>-1</sup>	1437 ± 2 cm <sup>-1</sup> 1438 ± 1 cm <sup>-1</sup>	1457 ± 3 cm <sup>-1</sup> 1459 ± 3 cm <sup>-1</sup>	1525 ± 5 cm <sup>-1</sup> 1523 ± 3 cm <sup>-1</sup>	1545 ± 5 cm <sup>-1</sup> 1545 ± 5 cm <sup>-1</sup>	1577 ± 3 cm <sup>-1</sup> 1577 ± 3 cm <sup>-1</sup>

frequency difference  $\Delta(v_{as} - v_s) \approx 60~\text{cm}^{-1}$  is predominant for the two salts at high concentrations, followed by the bridging bidentate coordination (III) (*Peak 2 – Peak 5*). The presence of both bridging and chelating complexes suggests that polymeric structures

(IV) involving several carboxylate headgroups also form at the surface.

Previous studies on specific interactions of lanthanum and yttrium ions with the carboxylate moiety in the bulk had sug-

**Table 2**Proposed assignments of the interfacial binding pattern for yttrium and lanthanum – carboxylate complexes.

$v_{\text{coo-,s}}$ - $v_{\text{coo-,as}}$	$\Delta(v_{as} - v_s)$	Coordination mode
Peak (1) – Peak (5)	$\approx 130 \text{ cm}^{-1}$	Ionic (ion pair)
Peak (2) – Peak (5)	$\approx 100 \text{ cm}^{-1}$	Bridging bidentate (III)
Peak (3) – Peak (4)	$\approx 60 \text{ cm}^{-1}$	Chelating bidentate (II)

gested similar types of complexes. For instance, a combined X-ray, IR, and NMR study on hydrated acetate crystals proposed bridging and polymeric complexes for La3+, and polymeric and chelating complexes for Y3+.[60] Another study using Raman spectroscopy targeting C—C stretching modes for determining the interactions of lanthanum with a hydrated lithium acetate, suggested the coexistence of both bridging and chelating bidentate structures, with a predominance of the chelating bidentate over the polymeric form. [61] However, many such studies in solid state or aqueous bulk solution are not directly transposable to the Langmuir monolayers. Given that the fatty acid monolayer is packed at the air/water interface, the possible carboxylate configurations and their respective interactions with trivalent cations are necessarily constrained. The stress added to the interfacial configuration (angles, distance, and strength of the chemical bonds) may induce variation in the frequency of vibration.

To confirm the assignments proposed for the planar geometry, a comparison is made with the lamellar bilayer structures of rare earth metallic soaps and their short chain homologues, for which single crystalline structures have been determined.[62-65] When the alkyl chain of these rare earth carboxylates is four carbons or longer, their configurations are similar to that of the all-trans hexagonal packing of the fatty acid in the UC phase. From the published crystalline structure of lanthanum butyrate monohydrate, [63] which is almost identical to that of the neodymium butyrate, [62] we examine in Fig. 5 the cross-section of the lamellar structure in the constrained environment of the monolayer. Both the chelating and bridging bidentate complexes can be readily identified (Fig. 5b), which is consistent with the assignments from the sum frequency spectra. The crystal structures from metallic soaps also provide an insight into the differences observed between La<sup>3+</sup> and Y<sup>3+</sup>. The yttrium cation has a coordination number of 8,[66] compared with 9 [63] for the lanthanum ion that is also 20% larger in size, [67] which provides the latter more flexibility to coordinate with neighbouring carboxylates via bridging complexes (note that a maximum of 6 carboxylate-oxygens can coordinate to the cations in the monolayer, with the remaining sites occupied by water-oxygens). Given that in Fig. 5 one of the butyrate lamellas was removed from the originally electrically neutral soap structure, the case shown would correspond to a charge overcompensated monolayer with a 2 (La<sup>3+</sup>): 3 (COO<sup>-</sup>) ratio. Overcharging is indeed observed for the highest salt concentrations, as discussed in the following section when evaluating the response from the OH stretching modes. However, at lower salt concentrations, when the monolayer charge is compensated in a 1:3 ratio, fewer ions will be present at the surface and the preferred binding complex will depend on the identity of the ion. A coarse analysis of how the trivalent ion will arrange at the surface for the charge compensated case in a packed carboxylic acid monolayer is presented in the SI (Figure S7).

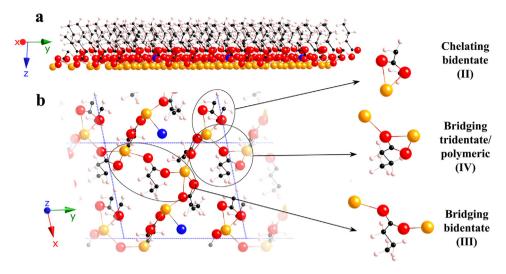
# 3.4. Targetting interfacial water molecules to detect the charge reversal of the surface

The OH stretching bands from interfacial water molecules can be sensitive probes for detecting changes in the surface charge. The bands report both on water molecules in direct proximity to the monolayer, as well as those found further into the bulk within the diffuse double layer that are perturbed by the decaying surface electric field. [68,69] The SSP polarized sum frequency spectra of arachidic acid monolayers on varying concentrations of YCl<sub>3</sub> in the subphase are presented in Fig. 6. Equivalent spectra for pure water and 1 mM NaCl are also added for reference.[46] The two broad bands centered at  $\sim 3250~{\rm cm^{-1}}$  and  $\sim 3500~{\rm cm^{-1}}$  on the pure water spectrum primarily originate from H<sub>2</sub>O molecules directly interacting with the uncharged carboxylic acid headgroup. [46,70] In the pure water case, contributions from molecules in the diffuse double layer are negligible (the Debye screening length in pure water is significantly longer than the non-linear coherence length of the sum frequency process, and most of the signal from the diffuse double layer cancels due to destructive interference [68,71]).

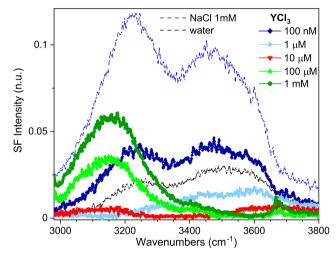
Upon addition of 100 nM YCl<sub>2</sub> to the subphase the OH bands increase relative to those of pure water (Fig. 6). The added intensity results primarily from water molecules in the diffuse layer, and is a direct indication of the onset of the monolayer deprotonation (Gouy-Chapman model predicts  $\sim$  3% deprotonation, see SI). The water bands are thus more sensitive probes than the carboxylate stretching modes to detect the first stages of charging (Fig. 2a, COO<sup>-</sup> features are below the detection limit at 100 nM). At 1 µM YCl<sub>3</sub>, when the monolayer is roughly 50% deprotonated (see C=O amplitude in Fig. 2c) and the Debye screening length is  $\sim$  2 times shorter, the OH contributions decrease below those from pure water. Had the surface charged been proportional to the amount of unprotonated carboxylic acid groups, the water signal would have been at least of the same order of magnitude as the one observed for 1 mM NaCl (Fig. 6). The discrepancy shows that yttrium cations not only promote deprotonation but upon adsorption and complexation with the carboxylate, they also reduce the net surface charge. By further increasing the salt concentration, the water bands reach close to zero intensity at 10 µM YCl<sub>3</sub>, before increasing again at higher concentrations, but with a broadband centered this time at  $\sim 3150~\text{cm}^{-1}$  and a sharper feature at  $\sim 3670 \text{ cm}^{-1}$ .

The trend involving a decrease of sum frequency intensity to a minimum, followed by an increase and a red-shift of the main contributing broad band, is indicative of a reversal of the surface charge, from negative to positive. This effect has been previously observed in several systems undergoing charge inversion, such as when adsorbing the positively charged cetyltrimethylammonium bromide surfactant on a negatively charged silica/water interface, [72] and at the  $TiO_2$ /water[73] and  $CaF_2$ /water[74,75] interfaces when varying the solution pH, among many other examples. [28,34] The charge reversal or overcompensation is typically driven by specific interactions of multivalent counterions with surfaces that display a relatively high charge density.[76-78] For the case of yttrium (and lanthanum), this strong interaction is demonstrated by the formation of binding complexes with the carboxylate headgroup, as revealed by vibrational sum frequency spectroscopy. The behaviour is also consistent with the changes observed in the surface pressure isotherms as a function of salt concentration (Fig. 1), where the most compressed monolayer is seen at 10  $\mu$ M YCl<sub>3</sub>, the concentration where charge neutralization is approximately reached.

The relatively sharp feature observed at  $\sim 3670~\text{cm}^{-1}$  provides additional support for the charge inversion at high YCl<sub>3</sub> concentrations. The high-frequency band is located in a similar frequency range as that observed for the OH stretches of Si-OH, Al-OH, and Ca-OH at the aqueous interface with silica,[42,79] alumina,[80] and calcium fluoride,[81] respectively. Thereby, it is assigned to the OH stretch of an yttrium hydroxide cation (i.e. Y(OH)<sub>2</sub><sup>+</sup> or Y (OH)<sup>2+</sup>). The change in sign of the surface potential due to charge overcompensation causes a significant increase in the concentra-



**Fig. 5.** Single monolayer extracted from the crystal structure of lanthanum(III) butyrate monohydrate presented along the (a) z-y plane and the (b) x-y plane. In the structure seen in the xy plane (b), several carboxylate coordination structures with the  $La^{3+}$  can be identified, *i.e.* chelating bidentate, bridging bidentate, and bridging tridentate (polymeric). They are consistent with the structures derived from the sum frequency spectra of arachidic acid Langmuir monolayers. In the figure, yellow spheres represent  $La^{3+}$  ions, red spheres carboxylate oxygens, and blue spheres selected water oxygens. The small black and white spheres are carbon and hydrogen atoms, respectively. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)



**Fig. 6.** Vibrational sum frequency spectra in the OH stretching region of perprotonated arachidic acid monolayers on solutions containing increasing concentrations of YCl<sub>3</sub>. Spectra were recorded under the SSP polarization combination at pH  $\sim$  5.8, 22 °C and a constant surface pressure of 20mN/m. Equivalent spectra on 1 mM NaCl and pure water are also added for reference. Spectra were recorded in triplicate, showing all the same trends.

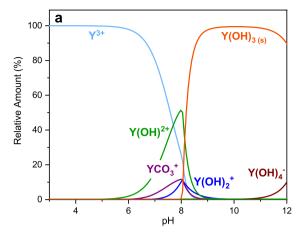
tion of surface anions, including  $OH^-$ , which can interact with the  $Y^{3+}$  ions already present at the surface. This will be further discussed in the next section.

# 3.5. Role of the surface charge on the surface speciation and detection limit

 $YCl_3$  as well as LaCl $_3$  show relatively intricate speciation diagrams at basic pH, with the formation of several hydroxide complexes, as shown in Fig. 7a for the yttrium case. Up to pH  $\sim$  7, the prevalent species is the free ionic form  $Y^{3+}$ . However, further increasing the pH shows the appearance of  $Y(OH)^{2+}$ ,  $Y(OH)^{2+}$ , and  $Y(OH)_{3(s)}$ , at the expense of the free ion. Conversely, below pH 6, when only the pH determining ions are present in solution, the

fatty acid monolayer is uncharged, but deprotonates in more alkaline conditions, reaching approximately 50% deprotonation at pH  $\sim$  10.8.[46] All sum frequency measurements discussed above were recorded at pH  $\sim$  5.8, where yttrium and lanthanum are in their free ionic form (i.e. Y³+ and La³+), and the fatty acid monolayer is at the deprotonation threshold (i.e., it can start to deprotonate even by the addition of small amounts of monovalent salts to solution).[46] To explore the influence of the bulk speciation and the intrinsic deprotonation state of the monolayer on the type of binding complexes formed, measurements were carried out at several different pH at a constant concentration of 1  $\mu$ M YCl₃. The corresponding sum frequency spectra recorded under the SSP polarization are shown in Fig. 7b.

The spectrum at pH  $\sim$  5.8 is similar to that presented in Fig. 2a for 1 µM YCl<sub>3</sub>, where the principal band in the symmetric carboxylate stretching region is centered at  $\sim 1458~\text{cm}^{-1}$ , and assigned to the chelating bidentate complex (marked as ii in Fig. 7b). The spectral features at the other pH values show, however, significant differences. At pH 5, when the yttrium ion is in its free ionic form Y<sup>3+</sup> but the monolayer does not as easily deprotonate, besides a weak broadband centered at  $\sim 1455~\rm cm^{-1}$  (ii), the spectrum shows a sharp contribution at  $\sim 1475~\rm cm^{-1}$  (i). This latter band was not observed at pH  $\sim$  5.8 at any salt concentration for either YCl<sub>3</sub> or LaCl<sub>3</sub>, and indicates that lowering the pH favours a different type of yttrium-carboxylate coordination pattern. When increasing the pH to 8, the main component in the symmetric carboxylate band is similar to that at pH  $\sim$  5.8 (ii), but with a much lower intensity. The formation of  $Y(OH)^{2+}$  and  $Y(OH)^{+}_{2}$  at the expense of  $Y^{3+}$  in the bulk, clearly has an influence on the behaviour observed as the hydroxides do not as readily interact with the carboxylate headgroup. At pH 9, there is only a very small indication of monolayer deprotonation, despite the fact that in the absence of trivalent ions the monolayer should be  $\sim 20\%$  deprotonated. [46] Given that at pH 9  $Y(OH)_{3(s)}$  is the main species and that almost all yttrium should precipitate out from solution, the effect observed must be caused by trace amounts of trivalent soluble species. Finally, at pH 11 when most yttrium is found in the trihydroxide form, a strong symmetric carboxylate stretch is seen at  $\sim 1408~\text{cm}^{-1}$ , which cor-



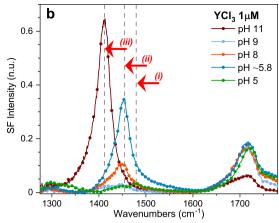
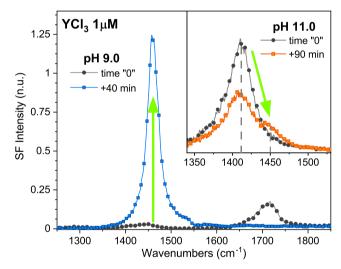


Fig. 7. a) Speciation diagram of YCl<sub>3</sub> at 25 °C calculated from HYDRA and MEDUSA software for a concentration of 1  $\mu$ M dissolved CO<sub>2</sub> (see SI for the full diagram). b) SSP polarized vibrational sum frequency spectra of deuterated arachidic acid on YCl<sub>3</sub> 1  $\mu$ M electrolyte solutions at pH 5, ~5.8, 8, 9 and 11 adjusted with either HCl, or NaOH. Temperature: 22oC. Surface pressure: 20mN/m. The spectra were collected 10 min after the deposition of the monolayer (time ="0"). Spectra were recorded in triplicate, showing all the same trends.

respond to the hydrated or solvent separated carboxylate, typically observed for monovalent ions.[26,46] Consequently, no traces of yttrium are detected with regards to the vibrational pattern of the carboxylate. However, the deprotonation fraction of the monolayer is lower than expected at this pH,[46] revealing residual metal-carboxylate interactions as developed below. All these results show that the speciation of yttrium in the bulk solution has, as expected, a strong influence on the observed behaviour. The formation of hydroxide complexes reduces the charge carried with the ionic yttrium species and the global amount of ionic species available, at the same time decreasing their propensity to migrate to the surface.

Time dependence: All data presented in Fig. 7b were recorded 10 min after the deposition of the monolayer on the electrolyte subphase (referred to as time "0"). However, in the pH range where the Y<sup>3+</sup> was not the dominant species, the spectra exhibited some time dependence. Fig. 8 shows spectra of the fatty acid monolayer at pH 9 and pH 11 collected at different times after deposition. Whereas at pH 9 almost no traces of deprotonation are detected at time = 0, the monolayer appears to almost fully deprotonate 40 min later, with a sharp peak centered at  $\sim 1459 \text{ cm}^{-1}$ , characteristic of the chelating binding pattern with yttrium ions. Similarly, at pH 11, the hydrated carboxylate contribution at  $1408 \text{ cm}^{-1}$  decreases at t = +90 min, while a shoulder at  $\sim 1458~\text{cm}^{-1}$  appears in the spectra. For both cases, the interaction of trivalent ions with the interface becomes obvious with time. It is worth noting that no such time dependence was observed in the spectra presented in the previous sections collected at pH  $\sim$  5.8.

The time dependence can originate from a diffusion limiting process. The charging of the fatty acid monolayer at high pH generates a negative surface potential that will attract cations to the interface. Although most of the yttrium species in the bulk at pH 9 and pH 11 should be in the hydrated hydroxide form  $Y(OH)_{3(s)}$ , sub-nanomolar concentrations of the ionic species  $Y(OH)^{2+}$ ,  $Y(OH)^{+}_{2}$ , and particularly  $Y^{3+}$ , remain in solution (see SI for an extended speciation diagram). Given the low concentrations, the diffusion of such species from the bulk solution to the interface will require some time. Moreover, since the hydronium ion concentration at the surface is a few orders of magnitude higher than in the bulk, [46,82] the speciation of yttrium in the surface region will also vary, with the preference to eventually form the free  $Y^{3+}$  ion. Con-



**Fig. 8.** Vibrational sum frequency spectra of deuterated arachidic acid on 1  $\mu$ M YCl<sub>3</sub> electrolyte solution at pH 9 and pH 11 (inset), directly after the deposition of the monolayer, and after a time of exposure (+40 min, +90 min). Temperature: 22 °C. Surface pressure: 20mN/m. Spectra were recorded in triplicate, showing all the same trends.

sequently, diffusion and surface speciation reaction kinetics will be time-limiting factors.

A time dependence behaviour was also observed in acidic conditions (i.e. pH = 5), when almost all yttrium in solution is found in its Y<sup>3+</sup> form, and ion diffusion to the surface is not the determining factor. In such a case, the underlying cause lies in the state of the monolayer deprotonation, as it is essentially fully uncharged in the absence of trivalent ions in solution (see SI, Figure S6). However, once the monolayer starts to deprotonate upon adsorption of the first ions, the charging accelerates the process. This effect can be used to improve the detection limits of trivalent ions in solution. In the presence of a background monovalent electrolyte (e.g. 100 mM NaCl), the monolayer can partially deprotonate at a lower pH when trivalent ions are more readily found in the free ionic form. For the case of yttrium, we show that the detection limit can be reduced to concentrations below 100 nM (see SI). These results also highlight the importance of carefully considering

the potential effects of trace amounts of trivalent ion impurities when studying charged interfaces [26].

### 4. Summary and concluding remarks

The carboxylate stretching modes were shown to be sensitive probes for identifying the different types of interactions between the trivalent ions  $La^{3+}$  and  $Y^{3+}$  and the carboxylic acid moiety. Vibrational sum frequency spectroscopy offers advantages when compared with the linear spectroscopies (e.g. IR and Raman), as the spectral features from both the symmetric and asymmetric carboxylate stretches are significantly better resolved. At least three different carboxylate species could be unambiguously identified when varying the salt concentration in solution. They were assigned to both ionic and covalent binding type interactions, [59] with the latter composed primarily of a mixture of bridging and chelating bidentate configurations. The classification relies on a generalised set of rules that have several exceptions.[15,16,18,58] It is hoped that the information provided here may stimulate additional theoretical, simulation-based, and complementary X-ray studies [57,83,84], to investigate further and confirm the binding structures formed by these trivalent ions at the liquid/vapour interface.

Although the same carboxylate species were identified for La $^{3+}$  and Y $^{3+}$ , they showed a different concentration dependence. In the case of yttrium, once the monolayer is completely deprotonated at  $\sim 5~\mu\text{M}$ , the relative proportions of the three different carboxylate species remain approximately constant. In contrast, for La $^{3+}$  the proportions continue to vary until 10 mM, where the chelating bidentate configuration reaches a maximum at the expense of the ionic contribution, which is no longer detectable at the highest salt concentrations. This difference is primarily attributed to the La $^{3+}$  ion size, which is almost 20% larger than Y $^{3+}$ [67], providing additional flexibility for finding the most energetically favourable binding pattern.

From the analysis of the spectral response of water molecules in the interfacial region, it was shown that specific interactions of Y $^{3+}$  with the carboxylic acid monolayer lead to a charge overcompensation at the surface, which classical electrostatic theories such as Gouy Chapman fail to predict. The charge reversal observed at concentrations above 10  $\mu M$  was further supported by the detection of a yttrium hydroxide complex at the surface.

Additionally, experiments performed at different subphase pH highlight the importance of the ion speciation in the surface region, which can be significantly different from that in the bulk, reflecting the expected changes in concentration of  ${\rm H_3O^+}$  or  ${\rm OH^-}$  close to the charged interface.

Finally, interactions between Y³+ and La³+ and the carboxylate headgroup could be detected at concentrations as low as 300 nM in solution. The detection limit was further reduced below 100 nM in the presence of a background electrolyte (i.e. NaCl) that promoted the partial deprotonation of the monolayer. It is anticipated that trivalent ions known to experience even stronger interactions with the carboxylate moiety, including, among others Co³+, Mn³+, Cr³+, Fe³+ or Sc³+,[85] could bring the detections limits down to sub nM concentrations. These results strongly indicate that trace amounts of trivalent ions in biological matrices and complex electrolytes could easily compete, if not dominate, the electrolyte interactions with the carboxylic acid moiety.

# **CRediT authorship contribution statement**

**Adrien Sthoer:** Conceptualization, Methodology, Validation, Formal analysis, Writing – original draft, Visualization. **Ellen M. Adams:** Investigation. **Sanghamitra Sengupta:** Investigation,

Validation. **Robert W. Corkery:** Methodology, Visualization. **Heather C. Allen:** Conceptualization, Supervision, Funding acquisition. **Eric C. Tyrode:** Conceptualization, Methodology, Validation, Formal analysis, Writing – review & editing, Supervision, Funding acquisition.

### **Declaration of Competing Interest**

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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### Appendix A. Supplementary material

Supplementary data to this article can be found online at <a href="https://doi.org/10.1016/j.jcis.2021.10.052">https://doi.org/10.1016/j.jcis.2021.10.052</a>. It includes: Langmuir isotherms of arachidic acid and deuterated arachidic acid on LaCl3-and YCl3 solutions at various concentrations and pH, details of the fitting procedure of the sum frequency spectra, predictions from the Gouy-Chapman model, estimation of proportions of metal-carboxylates contributions, sum frequency spectra showing improved detection limits in the presence of 100 mM NaCl background electrolyte, correlation between symmetric and asymmetric carboxylate stretches, detailed speciation diagrams for YCl3 and LaCl3, and coarse modelling of the trivalent ion adsorption on the fatty acid monolayer for the charge compensated case.

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